Doping effect of nano-diamond on superconductivity and flux pinning in MgB₂

C H Cheng¹, H Zhang^{1,2}, Y Zhao^{1,2,3}, Y Feng⁴, X F Rui², P Munroe¹, H M Zeng⁵, N Koshizuka⁶ and M Murakami⁶

- ¹ School of Materials Science and Engineering, University of New South Wales, Sydney 2052, NSW, Australia
- ² State Key Lab for Mesophysics, Department of Physics, Peking University, Beijing 100871, People's Republic of China
- ³ Superconductivity R&D Center, Mail Stop 152#, Southwest Jiaotong University, Chengdu, Sichuan 610031, People's Republic of China
- ⁴ Northwest Institute for Nonferrous Metal Research, PO Box 51, Xi'an, Shaanxi 710016, People's Republic of China
- ⁵ Key Laboratory of Polymeric Composites and Functional Materials, The Ministry of Education, Zhongshan University, Guangzhou 510275, People's Republic of China
- ⁶ Superconductivity Research Laboratory, ISTEC, 1-10-13 Shinonome, Koto-ku, Tokyo 135-0062, Japan

E-mail: c.cheng@unsw.edu.au

Received 30 January 2003, in final form 11 July 2003 Published DD MMM 2003 Online at stacks.iop.org/SUST/16/1

Abstract

Doping effect of diamond nanoparticles on the superconducting properties of MgB₂ bulk material has been studied. It is found that the superconducting transition temperature $T_{\rm c}$ of MgB₂ is suppressed by the diamond doping, however, the irreversibility field $H_{\rm irr}$ and the critical current density $J_{\rm c}$ are systematically enhanced. Microstructural analysis shows that the diamond-doped MgB₂ superconductor consists of tightly-packed MgB₂ nano-grains (\sim 50–100 nm) with highly-dispersed and uniformly-distributed diamond nanoparticles (\sim 10–20 nm) inside the grains. High density of dislocations and diamond nanoparticles may take the responsibility for the enhanced flux pinning in the diamond-doped MgB₂.

1. Introduction

Since the discovery of superconductivity at 39 K in MgB₂ [1], significant progress has been made in improving the performance of MgB₂ materials [2–7]. MgB₂ offers the possibility of wide engineering applications in the temperature range 20–30 K, where conventional superconductors, such as Nb₃Sn and Nb–Ti alloy, cannot play any roles due to their low T_c . However, the realization of large-scale applications for MgB₂-based superconductivity technology essentially relies on the improvement of the pinning behaviour of MgB₂ in high fields. As it has poor grain connection and a lack of pinning centres, MgB₂ often exhibits a rapid decrease in critical current density, J_c , in high magnetic fields. Fortunately, through the formation of nanoparticle structures in bulk

MgB₂ [3–5] and thin films [7], the problem of the poor grain connection can be solved, and the flux pinning force can also be significantly enhanced due to an increase of pinning centres served by grain boundaries. In order to improve further the performance of MgB₂, it is necessary to introduce more pinning centres, especially those consisting of nano-sized second-phase inclusions which often provide strong pinning forces.

Nano-diamond, prepared by the detonation technique, has been widely used as an additive to improve the performance of various materials [8]. However, nano-diamond-doping effect on the superconducting properties MgB_2 has never been reported although carbon with other forms has been used as dopants in MgB_2 [9–11]. The high dispersibility of the nano-diamond powder makes it possible to form a high

density of nano-inclusions in MgB_2 matrix. In this paper, we have investigated the doping effect of nano-diamond on the superconducting properties of MgB_2 . Our results show that the nano-diamond-doped MgB_2 consists of tightly-packed MgB_2 nano-grains ($\sim 50-100$ nm) with diamond nanoparticles ($\sim 10-20$ nm) wrapped within the grains. This unique microstructure provides the samples with a good grain connection for the MgB_2 phase and a high density of flux-pinning centres served by the diamond nanoparticles. Compared to the MgB_2 bulk materials doped with other nanoparticles [3–7], the irreversibility line has been significantly improved and the J_c in high magnetic fields has been largely increased in these nano-diamond-doped MgB_2 .

2. Experimental details

The MgB₂-diamond nano-composites with compositions of $MgB_{2-x}C_x$ (x = 0%, 5%, 8% and 10%) were prepared by solidstate reaction at ambient pressure. Mg powder (99% purity, 325 meshes), amorphous B powder (99% purity, submicron size) and nano-diamond powder (10-20 nm) were mixed and ground in air for 1 h. An extra 2% of Mg powder was added in the starting materials to compensate the loss of Mg caused by high-temperature evaporation. The mixed powders were pressed into pellets with dimensions of $20 \times 10 \times 3 \text{ mm}^3$ under a pressure of 800 kg cm⁻², sandwiched into two MgO plates, sintered in flowing Ar at 800 °C for 2 h, and then quenched to room temperature in air. In order to compare the substitution effect of carbon in boron in MgB2 with the additional effect of the nano-diamond in MgB2, a sample with an added 1.5 wt% of nano-diamond in MgB2 was prepared. The sintering temperature and the sintering time for this sample were reduced respectively to 730 °C and 30 min in order to reduce the chemical reaction between the MgB2 and the diamond. This sample has been referred to as '1.5 wt% C'.

The crystal structure was investigated by powder x-ray diffraction (XRD) using an X'pert MRD diffractometer with Cu $K\alpha$ radiation. The microstructure was analysed with a Philips CM200 field emission gun transmission electron microscope (FEGTEM). DC magnetization measurements were performed in a superconducting quantum interference device (SQUID, Quantum Design MPMS-7). J_c values were deduced from hysteresis loops using the Bean model. The values of the irreversibility field, H_{irr} , were determined from the closure of hysteresis loops with a criterion of 10^2 A cm⁻².

3. Results and discussion

Figure 1 shows the XRD patterns of the nano-diamond powder and the typical MgB₂-diamond composites. The reflection (111) of the diamond is extremely broad and an amorphous-phase-like background can be seen in the XRD pattern. The particle size of the nano-diamond powder is estimated to be about 20 nm according to the width of the reflection. In relation to the nano-diamond-doped MgB₂, one of the impurity phases is MgO, which may have formed during the mixing of raw materials in air. Diamond should be present as another impurity phase in the composites; however, its main reflection (111) cannot be seen in XRD patterns, due to an overlap with the MgB₂ (101) peak. With increasing

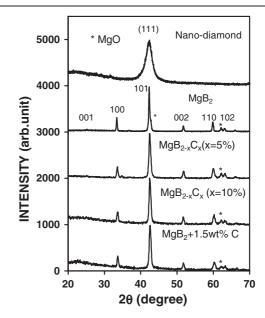


Figure 1. Powder XRD patterns for nano-diamond-doped MgB₂. The pattern on the top row is for the nano-diamond.

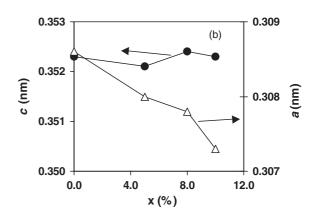


Figure 2. Doping dependence of lattice constants for nano-diamond-doped MgB₂.

doping level, an amorphous-phase-like background in the XRD pattern gradually appears, suggesting the existence of unreacted nano-diamond in the sample. As for the diamond-added MgB₂ sample (1.5 wt% C), which contains an x = 5.4% equivalent percentage of carbon atoms, the background of its XRD pattern shows some similarity to the background of the nano-diamond, suggesting that a substantial amount of unreacted nano-diamond exists within this sample.

Figure 2 shows the doping-level dependence of the lattice constants for the nano-diamond-doped MgB₂. The length of *c*-axis does not exhibit significant change with increasing nano-diamond doping level, but the *a*-axis is systematically decreased by doping the nano-diamond, indicating that a certain amount of carbon atoms have substituted for boron atoms in MgB₂. This result is consistent with those reported by other groups, which shows that partial substitution of boron by carbon results in a decrease of the lattice parameter [9–11].

Figure 3 shows the temperature dependence of magnetization at 2 mT for nano-diamond-doped MgB₂. The superconducting transition temperature $T_{\rm c}$ decreases with increasing doping level. This result is consistent with those reported by other groups on the carbon-doped MgB₂, which

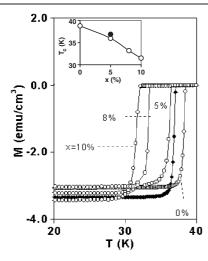


Figure 3. Temperature dependence of magnetization at 2 mT for nano-diamond-doped MgB_2 . Inset: doping-level dependence of T_c . The closed circles represent the results for the sample 1.5 wt% C in both the main and the inset figures.

suggests that carbon can partially substitute boron in MgB₂, decrease the carrier (hole) concentration and consequently decrease the T_c [9, 10]. The doping-level dependence of T_c is shown in the inset of figure 3. The values of onset T_c for these carbon-substituted MgB₂ samples are 38.6 K for x = 0%, 36.1 K for x = 5%, 33.0 K for x = 8% and 31.3 K for x = 10% (see the open circles). The T_c for the sample 1.5 wt% C is 36.9 K (see the closed circle), which is higher than that for the sample of x = 5% ($T_c = 36.1$ K), despite the former having a higher equivalent atomic percentage of carbon (x = 5.4%). This result suggests that only a part of carbon have been doped into the crystal structure of MgB₂ in the diamond-added sample, consistent with the XRD analysis.

Figure 4 shows the magnetic field dependence of J_c at 10, 20 and 30 K for the carbon-substituted MgB₂ samples. At 30 K, the undoped MgB₂ exhibits the highest J_c and the slowest decrease of J_c with H; whereas the sample of x=10% shows the lowest J_c and the quickest drop of J_c with H. It is evident that the J_c -H behaviour at 30 K for these samples is positively correlated to their T_c values. However, when the temperature decreases to the values far below T_c , a totally different situation appears. For example, at 10 and 20 K, the diamond-doped samples show a much better J_c -H behaviour. The J_c drops much more slowly in diamond-doped samples than in pure MgB₂. The best J_c at 20 K is found in the sample of x=10%, reaching a value of 6×10^3 A cm⁻² in a 4 T field, indicating that a strong flux pinning force exists in these diamond-doped samples.

However, the effect of diamond doping on the enhancement of flux pinning in MgB_2 may be counterbalanced by its suppression on superconductivity, as clearly shown in the situation of T=30 K (see figure 3). This counterbalancing effect may also exist at other temperatures, even when the effect of the J_c -enhancement is dominant. The further increase of J_c depends critically on reducing the T_c -suppression effect in the MgB_2 -diamond composite. This idea is confirmed by the results obtained in the diamond-added sample, 1.5 wt% C, which has a higher T_c than other diamond-doped samples (see figure 3) and contains more nano-diamond inclusions as suggested by the XRD analysis (see figure 1) and confirmed

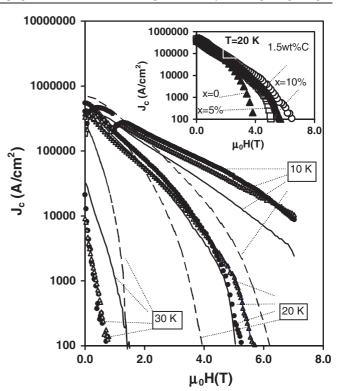


Figure 4. Magnetic field dependence of J_c at 10, 20 and 30 K for nano-diamond substituted MgB₂ with x = 0% (dashed lines), 5% (solid lines), 8% (solid circles) and 10% (opened triangles). Inset: comparison of J_c –H behaviour at 20 K between the diamond-added sample (1.5 wt% C) with some diamond-substituted samples (x = 0%, 5% and 10%).

by our TEM analysis shown below. As shown in the inset of figure 4, the diamond-added sample shows a much better J_c –H behaviour than the carbon-substituted sample. Its J_c reaches 1×10^4 A cm⁻² at 20 K and 4 T, and its $H_{\rm irr}$ reaches 6.4 T at 20 K. In fact, at all temperatures below 35 K, the J_c –H behaviour (results at 20 K are shown here only) of the diamond-added sample are much better than those of other samples in this study.

The $H_{\rm irr}-T$ relations for the diamond-doped MgB₂ are shown in figure 5. The $H_{\rm irr}$ (T) curves get steeper with increasing doping level, although the substitution of carbon for boron decreases the superconducting transition temperature. The best value of $H_{\rm irr}$ for these diamond-substituted MgB₂ samples including x=0%, 5%, 8% and 10% is found in the sample with x=10%, reaching 5.7 T at 20 K. However, the diamond-added MgB₂ (1.5 wt% C) shows a better $H_{\rm irr}-T$ behaviour than the diamond-substituted MgB₂. The $H_{\rm irr}$ value reaches 6.3 T at 20 K. The result clearly shows that the diamond doping does enhance the flux pinning in MgB₂ significantly, and also suggests that the pinning behaviour of the nano-diamond-doped MgB₂ depends on the density of the unreacted diamond nanoparticles in the samples.

Figure 6 shows the typical results from microstructural analysis for the diamond-substituted MgB₂ (figure 6(a)) and diamond-added MgB₂ samples (figure 6(b)). The diamond-substitutional sample mainly consists of relatively large MgB₂ grains ($\sim 1~\mu m$ or so in size) with a high density of dislocations. In some areas, discrete nano-sized particles can be seen. The diamond-added sample mainly consists of two kinds

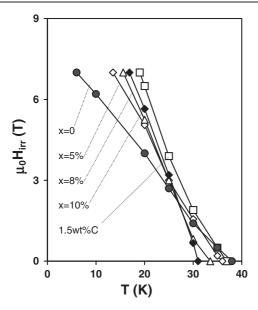


Figure 5. Variation of H_{irr} with temperature T for nano-diamond-doped MgB₂.

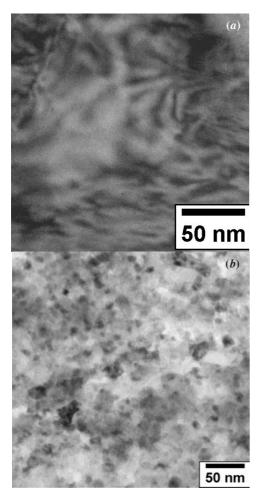


Figure 6. FEGTEM micrographs for (a) diamond-substituted MgB₂ with x = 5%; (b) diamond-added MgB₂ with the carbon content of 1.5 wt%. The atomic percentages of carbon in the sample 1.5 wt% C is equivalent to x = 5.4%.

of nanoparticles: MgB_2 grains with a size of 50–100 nm and diamond particles with a size of 10–20 nm. In fact, this diamond-added MgB_2 forms a typical nano-composite

material. The nano-diamond particles are inserted into the MgB₂ grains. As the *ab*-plane coherence length of MgB₂ is about 6–7 nm [12], these 10- to 20-nm-sized diamond inclusions, with a high density, are ideal flux pinning centres and are responsible for the high performance in our samples.

It is worth noting that the enhancement of flux pinning in nano-diamond-doped MgB2 is even better than the Ti-doped MgB₂ [3–5], where the TiB₂ nanoparticles mainly stay in grain boundaries, which results in a J_c of 9×10^4 A cm⁻² at 2 T and a H_{irr} of 4.6 T at 20 K (for a comparison, the corresponding values for the present system are J_c (20 K, 2 T) = 1.1 × 10^5 A cm⁻² and $H_{\rm irr}$ (20 K) = 6.3 T). This suggests that a highly-dispersed distribution of nano-inclusions in MgB₂ is more effective in enhancing flux pinning than a more localized distribution in the grain boundaries. Besides, compared to the Y_2O_3 nanoparticle doping which results in a J_c (20 K, 2 T) = $8 \times 10^4 \,\mathrm{A \ cm^{-2}}$ and H_{irr} (20 K) = 5.5 T, an advantage of the nano-diamond doping is that the lattice contact of the cubic diamond (a = 0.356 nm) is very close to the c-axis of MgB₂ (c = 0.352 nm). Therefore, these diamond nanoparticles may provide nucleation centres for MgB2 and are tightly bound to them. This may explain why our MgB2-diamond nano-composite performs much better than the nano-Y₂O₃doped MgB₂ [6]. It is expected that the performance of the MgB₂-diamond nano-composite may be further improved by optimizing the microstructure and the doping levels.

4. Summary and conclusions

In summary, we have investigated the doping effect of diamond nanoparticles on the superconducting properties of MgB_2 bulk material. We have observed that the superconducting transition temperature T_c of MgB_2 is suppressed by the diamond-doping, such as doping with other forms of carbon. However, the irreversibility field $H_{\rm irr}$ and the critical current density J_c are systematically enhanced by doping nano-diamond. Microstructural analysis shows that the diamond-doped MgB_2 superconductor consists of tightly-packed MgB_2 nano-grains ($\sim 50-100$ nm) with highly-dispersed and uniformly-distributed diamond nanoparticles ($\sim 10-20$ nm) inside the grains. High density of dislocations and diamond nanoparticles may be responsible for the enhancement of the flux pinning in the nano-diamond-doped MgB_2 .

Acknowledgments

The authors are grateful to Miss Sisi Zhao for her helpful discussion in preparing the manuscript. This work was supported in part by the University of New South Wales through the Goldstar Award for Cheng. Financial support from the Ministry of Science and Technology of China (NKBRSF-G19990646) is also acknowledged.

References

- [1] Nagamatsu J, Nakagawa N, Muranaka T, Zenitani Y and Akimitsu J 2001 *Nature* **410** 63
- [2] Cheng C H, Zhao Y, Feng Y, Zhu X T, Koshizuka N and Murakami M 2003 Supercond. Sci. Technol. 16 125

- [3] Zhao Y, Feng Y, Cheng C H, Zhou L, Wu Y, Machi T, Fudamoto Y, Koshizuka N and Murakami M 2001 Appl. Phys. Lett. 79 1154
- [4] Zhao Y, Huang D X, Feng Y, Cheng C H, Machi T, Koshizuka N and Murakami M 2002 Appl. Phys. Lett. 80 1640
- [5] Zhao Y, Feng Y, Machi T, Cheng C H, Huang D X, Fudamoto Y, Koshizuka N and Murakami M 2002 Europhys. Lett. 57 437
- [6] Wang J, Bugoslavsky Y, Berenov A, Cowey L, Caplin A D, Cohen L F, MacManus Driscoll J L, Cooley L D, Song X and Larbalestier D C 2002 Appl. Phys. Lett. 81 2026
- [7] Eorn C B et al 2001 Nature 411 558
- [8] Chen Q and Yun S 2000 Mater. Res. Bull. 35 1915
- [9] Takenobu T, Ito T, Chi D H, Prassides K and Iwasa Y 2001 Phys. Rev. B 64 134513
- [10] Mickelson W, Cumings J, Han W Q and Zettl A 2002 Phys. Rev. B 65 052505
- [11] Bharathi A, Jemima Balaselvi S, Kalavathi S, Reddy G L N, Sankara Sastry V, Hariharan Y and Radhakrishnan T S 2002 Physica C 370 211
- [12] Xu M, Kitazawa H, Takano Y, Ye J, Nishida K, Abe H, Matsushita A, Tsujii N and Kido. G 2001 Appl. Phys. Lett. 79 2779