# Improved irreversibility behavior and critical current density in MgB<sub>2</sub>-diamond nanocomposites

### Y. Zhao<sup>a)</sup> and C. H. Cheng

School of Materials Science and Engineering, University of New South Wales, Sydney 2052, NSW, Australia and State Key Lab for Mesophysics, Department of Physics, Peking University, Beijing 100871, China

## X. F. Rui and H. Zhang

State Key Lab for Mesophysics, Department of Physics, Peking University, Beijing 100871, China

#### P. Munroe

School of Materials Science and Engineering, University of New South Wales, Sydney 2052, NSW, Australia

#### H. M. Zeng

Key Laboratory of Polymeric Composites and Functional Materials, The Ministry of Education, Zhongshan University, Guangzhou, 510275, People's Republic of China

#### N. Koshizuka and M. Murakami

Superconductivity Research Laboratory, ISTEC, 1-10 13 Shinonome, Koto-ku, Tokyo, 135-0062, Japan

(Received 10 February 2003; accepted 8 July 2003)

MgB<sub>2</sub>-diamond nanocomposite superconductors have been synthesized by addition of nanodiamond powder. Microstructural analysis shows that the nanocomposite superconductor consists of tightly packed MgB<sub>2</sub> nanograins (~50–100 nm) with highly dispersed and uniformly distributed diamond nanoparticles (~10–20 nm) inside the grains. The  $J_c-H$  and  $H_{iir}-T$  characteristics have been significantly improved in this MgB<sub>2</sub>-diamond nanocomposite, compared to MgB<sub>2</sub> bulk materials prepared by other techniques. Also, the  $J_c$  value of  $1 \times 10^4$  A/cm<sup>2</sup> at 20 K and 4 T and the  $H_{irr}$  value of 6.4 T at 20 K have been achieved. © 2003 American Institute of Physics. [DOI: 10.1063/1.1606884]

Since the discovery of superconductivity at 39 K in MgB<sub>2</sub>,<sup>1</sup> significant progress has been made in improving the performance of  $MgB_2$  materials.<sup>2–4</sup>  $MgB_2$  offers the possibility of wide engineering applications in the temperature range 20-30 K, where conventional superconductors, such as Nb<sub>3</sub>Sn and Nb-Ti alloy, cannot play any roles due to their low  $T_c$ . However, the realization of large-scale applications for MgB<sub>2</sub>-based superconductivity technology essentially relies on the improvement of the pinning behavior of MgB<sub>2</sub> in high fields. As it has poor grain connection and a lack of pinning centers, MgB<sub>2</sub> often exhibits a rapid decrease in critical current density,  $J_c$ , in high magnetic fields. Fortunately, through the formation of nanoparticle structures in bulk  $MgB_2^{2-4}$  and thin films,<sup>5</sup> the problem of the poor grain connection can be solved, and the flux pinning force can also be significantly enhanced due to an increase of pinning centers served by grain boundaries. In order to improve further the performance of MgB<sub>2</sub>, it is necessary to introduce more pinning centers, especially those consisting of nanosized second-phase inclusions, which often provide strong pinning forces.

Nanodiamond, prepared by the detonation technique, has been widely used as an additive to improve the performance of various materials.<sup>6</sup> Yet, nanodiamond has never been used to increase the flux pinning force in MgB<sub>2</sub> superconductors until the present study. The high dispersibility of the nanodiamond powder makes it possible to form a high density of nanoinclusions in MgB<sub>2</sub> matrix. In this letter, we have prepared the MgB<sub>2</sub>-diamond nanocomposite, which consists of tightly packed MgB<sub>2</sub> nanograins (~50–100 nm) with diamond nanoparticles (~10–20 nm) wrapped within the grains. This unique microstructure provides the composite with a good grain connection for the MgB<sub>2</sub> phase and a high density of flux-pinning centers served by the diamond nanoparticles. Compared to the MgB<sub>2</sub> bulk materials prepared with other techniques, the irreversibility line has been significantly improved and the  $J_c$  in high magnetic fields has been largely increased in the MgB<sub>2</sub>-diamond nanocomposite.

The MgB<sub>2</sub>-diamond nanocomposites with compositions of MgB<sub>2-x</sub>C<sub>x</sub> (x=0%, 5%, 8%, and 10%) were prepared by solid-state reaction at ambient pressure. Mg powder (99% purity, 325 meshes), amorphous B powder (99% purity, submicron-size), and nanodiamond powder (10-20 nm) were mixed and ground in air for 1 h. An extra 2% of Mg powder was added in the starting materials to compensate the loss of Mg caused by high temperature evaporation. The mixed powders were pressed into pellets with dimensions of 20  $\times 10 \times 3$  mm<sup>3</sup> under a pressure of 800 kg/cm<sup>2</sup>, sandwiched into two MgO plates, sintered in flowing Ar at 800 °C for 2 h, and then quenched to room temperature in air. In order to compare the substitution effect of carbon in boron in MgB<sub>2</sub> with the additional effect of the nanodiamond in  $MgB_2$ , a sample with an added 1.5 wt% of nanodiamond in MgB<sub>2</sub> was prepared. The sintering temperature and the sintering time for this sample were reduced respectively to 730 °C and 30 min in order to reduce the chemical reaction between the

<sup>&</sup>lt;sup>a)</sup>Electronic mail: y.zhao@unsw.edu.au



FIG. 1. Powder XRD patterns for  $MgB_2$ -diamond nanocomposites. The pattern on the top row is for the nanodiamond.

 $MgB_2$  and the diamond. This sample has been referred to as 1.5 wt % C.

The crystal structure was investigated by powder x-ray diffraction (XRD) using an X'pert MRD diffractometer with Cu  $K\alpha$  radiation. The microstructure was analyzed with a Philips CM200 field emission gun transmission electron microscope (FEGTEM). Dc magnetization measurements were performed in a superconducting quantum interference device (Quantum Design MPMS-7).  $J_c$  values were deduced from hysteresis loops using the Bean model. The sample's dimensions with typical values of  $0.7 \times 2.1 \times 2.7$  mm<sup>3</sup> are used in the calculation of  $J_c$ . The values of the irreversibility field,  $H_{\rm irr}$ , were determined from the closure of hysteresis loops with a criterion of  $10^2$  A/cm<sup>2</sup>.

Figure 1 shows the XRD patterns of the nanodiamond powder and the typical MgB<sub>2</sub>-diamond composites. The reflection (111) of the diamond is extremely broad and an amorphous-phase-like background can be seen in the XRD pattern. The particle size of the nanodiamond powder is estimated to be about 20 nm according to the width of the reflection. In relation to the MgB<sub>2</sub>-diamond composites, one of the impurity phases is MgO, which may have formed during the mixing of raw materials in air. Diamond should be present as another impurity phase in the composites however, its main reflection (111) cannot be seen in XRD patterns, due to an overlap with the  $MgB_2$  (101) peak. As for the sample with the low doping level of x=5%, its XRD pattern looks the same as that of the undoped MgB<sub>2</sub>, except for a decrease of the lattice parameter along the a axis, indicating that a certain amount of carbon atoms have substituted for boron atoms in MgB<sub>2</sub>. This result is consistent with those reported by other groups, which show that partial substitution of boron by carbon results in a decrease of the lattice parameter.<sup>7,8</sup> With increasing doping level, an amorphous-phase-like background in the XRD pattern gradually appears, suggesting the existence of unreacted nanodiamond in the sample. As for the diamond added MgB<sub>2</sub> sample (1.5 wt % C), which con-

**PROOF COPY 004336A** 4% equivalent percentage of carbon atoms, the



FIG. 2. Magnetic field dependence of  $J_c$  at 10, 20, and 30 K for MgB<sub>2-x</sub>C<sub>x</sub> with x=0% (dashed lines), 5% (solid lines), 8% (solid circles), and 10% (opened triangles). Inset: superconducting transition curves for the diamond-doped samples. The closed circles represent the results for the sample 1.5 wt % C.

background of its XRD pattern shows some similarity to the background of the nanodiamond, suggesting that a substantial amount of unreacted nanodiamond exists within this sample.

The substitution of boron by carbon in our MgB<sub>2</sub> can also be reflected by the gradual decrease of  $T_c$  with increasing carbon content (see the inset of Fig. 2). The values of onset  $T_c$  for these carbon-substituted MgB<sub>2</sub> samples are 38.6 K for x=0%, 36.1 K for x=5%, 33.0 K for x=8%, and 31.3 K for x=10%. The  $T_c$  for the sample 1.5 wt % C is 36.9 K, which is higher than that for the sample of x=5% ( $T_c$ = 36.1 K), despite the former having a higher equivalent atomic percentage of carbon (x=5.4%).

Figure 2 shows the magnetic field dependence of  $J_c$  at 10, 20, and 30 K for the carbon-substituted MgB<sub>2</sub> samples. At 30 K, the undoped MgB<sub>2</sub> exhibits the highest  $J_c$  and the slowest decrease of  $J_c$  with H; whereas the sample of x=10% shows the lowest  $J_c$  and the quickest drop of  $J_c$  with H. It is evident that the  $J_c-H$  behavior at 30 K for these samples is positively correlated to their  $T_c$  values. However, when the temperature decreases to the values far below  $T_c$ , a totally different situation appears. For example, at 10 and 20 K, the diamond-doped samples show a much better  $J_c-H$  behavior. The  $J_c$  drops much more slowly in diamond-doped samples than in pure MgB<sub>2</sub>. The best  $J_c$  at 20 K is found in the sample of x=10%, reaching a value of  $6 \times 10^3$  A/cm<sup>2</sup> in a 4 T field, indicating that a strong flux pinning force exists in these diamond doped samples.

The  $H_{irr}-T$  relations for the diamond-substituted MgB<sub>2</sub> are shown in the inset of Fig. 3. The  $H_{irr}(T)$  curves get steeper with increasing doping level. The best value of  $H_{irr}$ 



FIG. 3. Variation of  $H_{irr}$  with reduced temperature  $T/T_c$  for MgB<sub>2-x</sub>C<sub>x</sub> with x=0%, 5%, 8%, and 10%. Inset:  $H_{irr}-T$  plot for the same data shown in the main figure.

reaches 5.7 T at 20 K for the sample of x=10%. As the  $T_c$  values vary with the diamond-doping level, only the  $H_{irr}-T$  relation cannot directly reflect the intrinsic irreversibility behavior for the samples of different doping levels. In the main panel of Fig. 3, the temperature dependence of  $H_{irr}$  is replotted using a reduced temperature,  $T/T_c$ . It is evident that the irreversibility field shifts towards higher temperatures with the increase of the diamond-doping level. The result clearly shows that the diamond doping does enhance the flux pinning in MgB<sub>2</sub> significantly.

However, the effect of diamond doping on the enhancement of flux pinning in MgB<sub>2</sub> may be counterbalanced by its suppression on superconductivity, as clearly shown in the situation of T=30 K (see Fig. 2). This counterbalancing ef-



FIG. 4. Comparison of  $J_c-H$  relations at 20 K for diamond-added MgB<sub>2</sub> sample 1.5 wt% C with diamond-substituted MgB<sub>2</sub>. The atomic percentages of carbon in the sample 1.5 wt% C and the sample of x=5% are almost the same. Inset:  $H_{\rm irr}-T$  relations for the same samples shown in the

fect may also exist at other temperatures, even when the effect of the  $J_c$  enhancement is dominant. The further increase of  $J_c$  depends critically on reducing the  $T_c$ -suppression effect in the MgB<sub>2</sub>-diamond composite. This idea is confirmed by the results obtained in the diamondadded sample, 1.5 wt % C, which has a higher  $T_c$  than other diamond-doped samples (see inset of Fig. 2) and contains more nanodiamond inclusions as suggested by the XRD analysis (see Fig. 1) and confirmed by our transmission electron microscopy analysis shown later. As shown in Fig. 4, the diamond-added sample shows a much better  $J_c-H$  behavior than the carbon-substituted sample. Its  $J_c$  reaches 1  $\times 10^4$  A/cm<sup>2</sup> at 20 K and 4 T, and its H<sub>irr</sub> reach 6.4 T at 20 K. In fact, at all temperatures below 35 K, the  $J_c-H$  behavior (results at 20 K are shown here only) and the  $H_{irr}-T$ relation (see the inset of Fig. 4) of the diamond-added sample are much better than those of other samples in this study.

Figure 5 shows the typical results from microstructural analysis for the diamond-substituted  $MgB_2$  and diamond-



FIG. 5. FEGTEM micrographs for (a) a typical grain ( $\sim 1 \ \mu$ m) of diamond substituted MgB<sub>2</sub> with x=5% which shows the high density of dislocations (dark stripes) in the sample; (b) an enlarged view of the dislocations in Fig. 5(a); (c) diamond added MgB<sub>2</sub> with the carbon content of 1.5 wt %. The grain boundaries of MgB<sub>2</sub> are indicated by the guidelines. The diamond nanoparticles are marked by letter A beside it (for small ones) or on the particles (for big ones). The atomic percentages of carbon in these two samples are almost the same.

#### 4 Appl. Phys. Lett., Vol. 83, No. 10, 8 September 2003

added MgB<sub>2</sub> samples. The diamond-substitutional sample mainly consists of relatively large MgB<sub>2</sub> grains (~1  $\mu$ m or so in size) with a high density of dislocations. In some areas, discrete nanosized particles can be seen [Figs. 5(a) and 5(b)]. The diamond-added sample mainly consists of two kinds of nanoparticles: MgB<sub>2</sub> grains with a size of 50–100 nm and diamond particles with a size of 10–20 nm [see Fig. 5(c)]. In fact, this diamond-added MgB<sub>2</sub> forms a typical nanocomposite material. The nanodiamond particles are inserted into the MgB<sub>2</sub> grains. As the *ab* plane coherence length of MgB<sub>2</sub> is about 6–7 nm,<sup>9</sup> these 10–20 nm-sized diamond inclusions, with a high density, are ideal flux pinning centres and are responsible for the high performance in our samples.

The significant improvement of  $J_c$  and  $H_{irr}$  in the nanodiamond-added samples (1.5 wt%) can be attributed to their nanocomposite structure which consists of two kinds of nanoparticles: MgB<sub>2</sub> grains with a size of 50-100 nm and diamond particles with a size of 10-20 nm. The enhanced number of grain boundaries associated with the smaller grain size can enhance the flux pinning, as reported previously.<sup>2-5</sup> However, only this factor cannot fully explain the experimental results because the enhancement of flux pinning in the nanodiamond-added samples (1.5 wt%) is even much better than that in the Ti-doped  $MgB_2^{2-4}$  where the average grain size of MgB<sub>2</sub> reaches 8-10 nm. This indicates that there may be other mechanisms of flux pinning enhancement in the present system. One of the most likely candidates is the diamond nanoparticles which may play a similar role as  $Y_2O_3$  nanoparticles did in  $Y_2O_3$ -doped MgB<sub>2</sub>.<sup>10</sup> It is worth noting that, compared to the  $Y_2O_3$   $Y_2O_3$ -doped MgB<sub>2</sub> the nanodiamond-added samples (1.5 wt %) has a higher  $J_c$  and  $H_{\rm irr}$ . This may be due to the advantage of the nanodiamond whose lattice contact (for the cubic diamond a = 0.356 nm) is very close to the c axis of MgB<sub>2</sub> (c = 0.352 nm). Therefore, these diamond nanoparticles may provide nucleation centres for MgB<sub>2</sub> and are tightly bound to them. It has been reported that some undoped MgB<sub>2</sub> samples with a slight grain texture also show a high  $J_c$  and  $H_{irr}$  (see, for example, Narozhnyi et al.),<sup>11</sup> suggesting that achieving a textured microstructure is another effective way to improve  $J_c$  of MgB<sub>2</sub> because of a slight anisotropy existing in this system. Accordingly, it is expected that the performance of the MgB<sub>2</sub>-diamond nanocomposite may be further improved by optimizing the microstructure and the doping levels.

In summary, we have synthesized a MgB<sub>2</sub>-diamond nanocomposite superconductor by adding nanodiamond powder into MgB<sub>2</sub>. The nanocomposite consists of tightly packed MgB<sub>2</sub> nanograins (~50–100 nm) with diamond nanoparticles (~10–20 nm) inserted inside these grains. The  $J_c-H$  and  $H_{irr}-T$  characteristics have been significantly improved in this MgB<sub>2</sub>-diamond nanocomposite, in comparison with MgB<sub>2</sub> bulk materials prepared with other techniques.

The authors are grateful to Sisi Zhao for her helpful discussion in preparing the manuscript. This work was supported in part by the University of New South Wales (Goldstar Award for Cheng). Financial support from the Ministry of Science and Technology of China (NKBRSF-G19990646) is also acknowledged.

- <sup>1</sup>J. Nagamatsu, N. Nakagawa, T. Muranaka, Y. Zenitani, and J. Akimitsu, Nature (London) **410**, 63 (2001).
- <sup>2</sup>Y. Zhao, Y. Feng, C. H. Cheng, L. Zhou, Y. Wu, T. Machi, Y. Fudamoto,
- N. Koshizuka, and M. Murakami, Appl. Phys. Lett. 79, 1154 (2001).
- <sup>3</sup>Y. Zhao, D. X. Huang, Y. Feng, C. H. Cheng, T. Machi, N. Koshizuka, and M. Murakami, Appl. Phys. Lett. 80, 1640 (2002).
- <sup>4</sup>Y. Zhao, Y. Feng, T. Machi, C. H. Cheng, D. X. Huang, Y. Fudamoto, N. Koshizuka, and M. Murakami, Europhys. Lett. **57**, 437 (2002).
- <sup>5</sup>C. B. Eorn, M. K. Lee, J. H. Choi, L. J. Belenky, X. Song, L. D. Cooley,
- M. T. Maus, S. Patnaik, J. Jiang, M. Rikel, A. Polyanskii, A. Gurevich, X.
- Y. Cai, S. D. Bu, S. E. Babcock, E. E. Hellstrom, D. C. Larbalestier, N.
- Rogado, K. A. Regan, M. A. Hayward, T. He, J. S. Slusky, K. Inumaru, M.
- K. Haas, and R. J. Cava, Nature (London) 411, 558 (2001).
- <sup>6</sup>Q. Chen and S. Yun, Mater. Res. Bull. **35**, 1915 (2000).
- <sup>7</sup>T. Takenobu, T. Ito, D. H. Chi, K. Prassides, and Y. Iwasa, Phys. Rev. B **64**, 134513 (2001).
- <sup>8</sup>W. Mickelson, J. Cumings, W. Q. Han, and A. Zettl, Phys. Rev. B 65, 052505 (2002).
- <sup>9</sup>M. Xu, H. Kitazawa, Y. Takano, J. Ye, K. Nishida, H. Abe, A. Matsushita, N. Tsujii, and G. Kido, Appl. Phys. Lett. **79**, 2779 (2001).
- <sup>10</sup>J. Wang, Y. Bugoslavsky, A. Berenov, L. Cowey, A. D. Caplin, L. F. Cohen, J. L. MacManus Driscoll, L. D. Cooley, X. Song, and D. C. Larbalestier, Appl. Phys. Lett. **81**, 2026 (2002).
- <sup>11</sup> V. N. Narozhnyi, G. Fucks, A. Handstein, A. Gumbel, J. Eckert, K. Nenkov, D. Hinz, O. Gutfleisch, A. Walte, L. N. Bogacheva, I. E. Kostyleva, K.-H. Muller, and L. Schultz, cond-mat/0206513.

Yo,